

AE3030 Notes up till page 75

I'm going to type up everything I know about Aerodynamics SO FAR
It is September 10th today. The first quiz is on September 22nd

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PROPERTIES OF FLUIDS

definite mass but indefinite volume

viscosity

--solids undergo static deformation

--fluid is a substance that deforms continuously under the action of a shearing force

--if there is zero shear stress, then we are in **hydrostatic stress condition**

CONTINUUM

-avg behavior of particles/molecules

--1 mole of O₂ is 6.02×10^{23} O₂'s

-Mean free path

The avg distance a particle travels before colliding with another

Knudsen number

-This is the mean free path divided by a characteristic length scale, such as chord length or diameter or airfoil length

When the Knudsen number is small, the continuum model is of value

Kn number is represented by λ

If the Knudsen number is 0.1 and under, we can assume continuum

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Some dimensional analysis

1 slug is $1 \text{ lbf s}^2 / \text{ft} = 1 \text{ lbm}(32.2 \text{ ft/s}^2)$

$1 \text{ lbf} = 1 \text{ lbm} * 32.2 \text{ ft/s}^2$

FLUID UNITS

Density

is mass over volume

m / l^3

kg / m^3

slug/ft³

Incompressible flows when changes in density can be neglected -- no fluid is ever actually incompressible

-formal definition

density = the limit as the change in volume approaches zero of the change in mass over the change in volume

Density is a thermodynamic property, meaning it relates to Temperature AND pressure

Density is also equal to P/RT (the R is the specific gas constant)

$PV = nRT$ (the universal gas law)

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Pressure

Pressure is the rate of change of momentum (?)

Pressure = Force / Area = Pa = N/m² = lbf/ft²

Make sure you know that $P - P_{\text{inf}}$ is gauge pressure, which is comparing absolute pressure against a reference pressure. There is no such thing as a negative absolute pressure.

#####

Temperature

A particle by itself has 3 degrees of translational motion and a group of particles (a duo of particles) has 2 degrees of rotational motion (the rotation about its axis is ignored because it's so small)

Altogether that makes 5 degrees of freedom

Calorically perfect gas

: c_p and c_v are constant. $c_v = 5/2 * R$

A thermally perfect gas is when c_v and c_p are functions of temperature

Real gas is when c_v and c_p are functions of temperature and pressure

UNITS

K and R

K is C + 273.15

FLOW VELOCITY

-u,v,w: bulk velocities / average velocity

CONSERVATION EQUATIONS

Definition of a streamline

-a curve whose tangent at any point is in the direction of the velocity vector at that point

t = constant

streamlines in 3D are stream TUBES

Unsteady Flow - when the streamlines change

In order to use bernoulli's eq, you need to make sure the flow is steady, so make sure to pick an appropriate reference frame!

Conservation of Mass/ Continuity Equations

The integral with respect to time over a control volume of density times the volume differential (unsteady term) plus the line integral over a control surface of the density multiplied with the velocity dotted with a normal vector multiplied by that control area (flux term) differential is all equal to zero.

$$d/dt \int_{cv} (\rho \cdot dV) + \text{line-int}_{cs} (\rho \cdot (\mathbf{v} \cdot \mathbf{n}) \cdot dA) = 0$$

When you assume steady flow, you can ignore the first term.

$-\rho_1 \cdot v_1 \cdot A_1$ is the mass flow rate in

$\rho_2 v_2 A_2$ is the mass flow rate out

If you assume compressible, the densities cancel

$$v_1 A_1 = v_2 A_2$$

Conservation of Momentum Equation

Rate of change of momentum is equal to the mass flow rate times velocity

(derivations on page 6 and 7)

Euler's equation:

$$\rho \cdot V \cdot dV + dP = 0$$

(assumptions: steady, inviscid, no body forces)

Bernoulli's equation:

$$P_o = P + \rho \cdot v^2 / 2$$

assumptions: incompressible, steady, inviscid, no body forces

Thermodynamics 1st law

$$\Delta U = Q - W$$

Internal energy change is equal to the heat added to the system minus the work done by the system

$$U = cv \cdot T$$

Other forms of energy

-Kinetic energy

-Potential energy (gravitational, electric potential) -- largely ignored

-Assumptions: steady, no body forces, inviscid flow, adiabatic

Work done - viscous force (ignore)

pressure flow work

-shaft work (ignore)

Rate at which work is being done is called Power

$$\text{Power} = F \cdot \text{velocity}$$

(derivation on page 9-10)

Energy equation derivation continues on page 12

$P = \rho \cdot R \cdot T$ is called an equation of state

Enthalpy and relations between R, Cv, Cp, and gamma (maybe)

Stagnation enthalpy

$$h_0 = h + \frac{1}{2}v^2$$

Static enthalpy is equal to $c_p \cdot T$

c_p is equal to $c_v + R$

Summary of Energy

for a stream tube:

$-\rho \cdot A \cdot V$ is constant

$pVdV + dP = 0 \Rightarrow$ Euler's equation

$h + \frac{v^2}{2} = h_0 = \text{constant}$

ASSUMING: steady, inviscid, no body forces, adiabatic

2nd LAW OF THERMODYNAMICS

$ds \geq \delta q / T$ $ds = \delta q / T + ds_{\text{irr}}$ and $ds_{\text{irr}} \geq 0$, so everything is irreversible

The total entropy of a system and surroundings is always increasing or at best stays the same

entropy \Rightarrow state variable (path independent)

heat added \Rightarrow path dependent

ds you can integrate

δq you cannot integrate, hence the weird lower-case delta

Define a Reversible process:

-one in which no dissipative phenomenon occurs

i.e., viscosity, thermal conductivity & mass diffusion are absent

-isentropic flow \Rightarrow adiabatic and reversible

Non-isentropic

\Rightarrow shock so that P_0 decreases

\Rightarrow boundary layers / viscous fluids

One form of the 2nd law:

$$s_2 - s_1 = c_p \ln(T_2/T_1) - R \ln(P_2/P_1)$$

Isentropic means that $s_2 - s_1 = 0$

you can derive a temperature – pressure relation for isentropic flow (derivation on page 14)

Also, there is another version of the 2nd law

$$s_2 - s_1 = c_v \ln(T_2/T_1) + R \ln(v_2/v_1) \text{ where } v \text{ is specific volume in this case}$$

INTRO TO AIRFOILS (page 15)

uses of airfoils:

-wings, props and fans, helicopters, compressors and turbines, hydrofoils, wind turbines, etc

Evolution of airfoils.

in 1873 they were like a line,... had a sharp leading edge, in 1903 there was a thicker leading edge but the whole thing was still thin – biplane

by 1917 we got the modern airfoil

AIRFOIL VOCAB

Chordline

-connects leading edge and trailing edge

camber

-is in terms of the chordline

thickness

– perpendicular to camber line, midpoint is on camber line

Symmetric airfoil

→ camber collapses to the chordline

angle of attack – angle between freestream and the chordline

Units for angle of attack

→ always use radians!!

alpha is not constant across a wing – wings are twisted --- rotor blades, too, but even moreso
The root of the wing typically has a higher alpha and the tip has lower, and maybe even negative alpha

FORCES

Lift

-force acting perpendicular to V_{∞}

Drag

-Force acting along as V_{∞}

Moment b/c the l and d are distributed along the foil... convention is that clockwise is positive

MORE FORCES

pressure distribution P

units in Pascals

shear stress distribution

units in Pa Force / length²

Integrate these distributions to obtain the resultant Lift, Drag, and Moments

NORMAL AND AXIAL FORCES

They are with respect to the chord line as opposed to the V_{∞} free stream

L and D are dimensional

L' and D' are per unit span (like airfoil)

But we need to non-dimensionalize them

Lift has units of N

So we need an N to help

we got $\frac{1}{2} \rho_{\infty} v_{\infty}^2 s$ (s is chord times span for wings) (this term is dynamic pressure q_{∞})

COEFFICIENTS

Cl with lowercase is for a 2D airfoil

CL is $L / (q_{\infty}) * s$ (area)

Cl is $L' / (q_{\infty}) * c$ (chord)

So.....

$$L = q_{\infty} s c_L$$

Similarly

$$C_d = D' / q_{\infty} c$$

etc etc

There's also moment

$$cM = M' / (q_{\infty}) c^2 \quad (\text{it's } c^2 \text{ 'cause units})$$

$$c_l, c_d, c_m = f(\text{geometry, } \alpha, \text{Mach number, Reynold's number})$$

Center of Pressure

It is the location where the resultant of the distributed load effectively acts on the body – if moments are taken about the CP, the integrated effect of the distributed load would be zero

Pressure Coefficient - C_p

(graph on page 22)

On a graph of C_l vs. α , slope is usually 2π

For a symmetric airfoil, when $\alpha = 0$, C_l is also zero.

$$C_l = 2\pi(\alpha - \alpha_0) \quad (\text{valid for the inviscid regime of all airfoil, subsonic})$$

What about DRAG?

There are three sources of drag:

- Pressure
- Viscous
- Wave

Pressure drag is caused by there being high pressure on the leading and trailing edge and there being low pressure on the top and bottom. There is an adverse pressure gradient when the air flows from high p at the TL to the low p in the middle area

Pressure drag depends on shape

Viscous drag

Occurs at the surface

-liquid wants to bring stuff with it

Wave drag comes from shocks and it's only for transonic and supersonic

WHAT IS VISCOSITY?

μ is the coefficient of viscosity (dynamic)

μ is in units of Newton-seconds per meter²

OR pascal-seconds

OR

kg / meter-second

μ is a function of temperature. At a certain temperature, the μ for air is about $1.983 \cdot 10^{-5}$ Pa-s

$\mu_{\text{water}} = 1.00 \cdot 10^{-3}$ Pa-s

$\mu_{\text{liquid honey}} = 10$ Pa-s

Remember, viscosity is a property of the fluid, and not the flow!!!!

.....

SIMILARITY PARAMETERS

Reynold's number (Re)

$Re = \rho \cdot V \cdot L / \mu$

density * velocity * Length reference / dynamic viscosity

$Re_{\text{inf}} = \rho_{\text{inf}} \cdot u_{\text{inf}} \cdot L / \mu_{\text{inf}}$

OR

you can have different Re at different points

-Mach Number: ratio of flow velocity to speed of sound

$$M = V / a$$

Local velocity/local speed of sound

$$a = \sqrt{\gamma * R * T}$$

VISCOUS DRAG

When fluid is very close to the surface, its velocity is 0

In a very short distance, fluid speeds up to its full potential.
This distance is called the boundary layer.

Shear stress at a wall

$$\tau_{\text{wall}} = \mu * dv/dn \text{ for wall}$$

LAMINAR FLOW

Lamina (a thin scale or sheet)

in laminar flow, molecular diffusion is the only mechanism for energy transfer, so it's slow

Turbulent flow is chaotic, irregular, 3d, disorderly and random, unsteady, and you have eddies

Mixing is from eddies

For turbulent flow, the slope dv/dn is high
The air speeds up faster because it's mixing better

$$\tau_{\text{wall}} = \mu * dv/dn \text{ so}$$

$\tau_{\text{wall, turb}}$ is greater than $\tau_{\text{wall, laminar}}$

And

skin friction is higher during turbulent flow

Wince tau is tangential to the airfoil, we can define an x_component for drag as tau_x, which is tau_wall*dx

D'_viscous = the integral of tau_lower*dx PLUS tau_upper*dx

You can non-dimensionalize all this by dividing each term by q_inf (dynamic pressure) to get the cf, which is the SKIN FRICTION COEFFICIENT

cf is a function of position

cd, viscous is an INTEGRATED QUANTITY

TRANSITION

This is when laminar flow turns to turbulent flow and it's very different from serpartation

something going to speed up the transition:

- rough surface
- more turbulence in the free stream
- adverse pressure gradient
- heating the fluid surface

An adverse pressure gradient means that the pressure increases along the chord length
 $bP/bx > 0$ in other words.

SEPARATION

(there is a good drawing page 29, bottom)

Separation occurs when $bv/bn = 0!!!$

Turbulent boundary layers have more energy close to the surface, and therefore it's harder for the boundary layer to separate.

(good pictures, explanation page 30)

STREAMLINES, PATHLINES, STEADY vs. UNSTEADY, ETC

Equation of a streamline is....

(derivation on page 31)

basically it is velocity vector cross with the path vector, and that equals zero

You get $dx/u = dy/v = dz/w$

PATHLINE

Lie a long exposure shot – tracks particles
you integrate the velocity

$dx/dt = u(x,y,z,t)$ etc

IN A STEADY FLOW

The pathline = the streamline

ANGULAR VELOCITY

-fluids rotate and deform at the same time, so we do perpendicular lines and average out

(The picture with the description and derivation is on page 32)

(A good picture of the C_p chart with pressure distribution is on page 41)

For the upper surface, at x/c is zero, you get the stagnation point, which has the highest pressure.

At the peak of the upper surface, you get the absolute lowest pressure and the highest flow velocity

On the back end you have an adverse pressure gradient (caused by separation)

Highest $C_p = 1$ at incompressible flows

if $C_p > 1$, your flow is compressible!!!

Integrate the C_p vs x/c curve to get the coefficient of lift!!

ANGULAR VELOCITY

We get another derivation thing on page 42, 43

BASICALLY Angular velocity is rate of change of angular rotation $d\theta_1/dt$ and $d\theta_2/dt$.

$d\theta_1/dt$ is equal to the limit as t goes to zero of $-b_u/b_y$

$d\theta_2/dt$ is the limit as t approaches zero of b_v/b_x

ANGULAR VELOCITY - $w_z = \frac{1}{2} (d\theta_1/dt + d\theta_2/dt)$ (average) = $\frac{1}{2} (bv/bx - bu/by)$

For 3 dimensions.....

formula on page 43 (it'll look terrible if I write it out here)

What is vorticity?>

Vorticity is defined by TWICE the angular velocity

Denoted by squiggle...it is a vector

The curl of something is the gradient crossed with it (upside-down triangle)

so the curl of velocity is

gradient X velocity

in 3-dimensions

$(b/bx_i + b/by_j + b/bz_k) \times (x_i + y_j + z_k)$

If the curl is equal to zero, then there is no rotation.

If the curl; is not equal to zero, then the flow is rotational and has rotational velocity!

RETURNING TO FLUID ELEMENTS

defining the inside angle of the deforming fluid element as kappa, or K

K is equal to $\theta_1 - \theta_2$

$-\Delta K$ is equal to strain that's positive (look at page 45 for details)

"Dilatation" of a fluid is equal to the divergence of the fluid

gradient DOT velocity

which is

$bu/bx + bv/by + bw/bz$

REMINDER: ABSOLUTE PRESSURE CANNOT BE NEGATIVE

KNOWLEDGE OF VISCOUS FLOW IS ESSENTIAL FOR EXTERNAL FLOW APPLICATIONS

-airfoil wing design

-airfoil and wing analysis

-fans, props, helicopters

-wind turbines

-cars, ship and sail shapes

- channels, ducts, and pipes
- transport
- biomedical applications
- heart pumps

flow through arteries, valves

Viscous flow applications – inacting flow

- mizing of species obeys a set of governing laws and equations that are similar to the viscous transport of momentum and energy from one region of the fluid to another

turbulent flow has much more surface area!!

Properties of flow....

viscosity μ

is a property of the fluid no the flow AGAIN

it's derived from intermolecular forces

- normal vs. tangential stress

Viscosity in liquids

As Temperature increases, μ decreases b/c liquid expands, molecules move apart

in GASSES, viscosity is associated with random molecule movement, so as temperature increases, μ increases

FLUIDS

Reminder: strain is infinite, HOOKE'S LAW does not apply

Therefore we should define rate of straing

Newton

- tangential stress is proportional to rate of strain

$\tau = \mu * \text{velocity} / \text{height}$

For his weird plate project

NEWTON'S ESTIMATES OF VISCOSITY

$\tau = F/A = \mu * V/h$

This only works for one dimension

STOKES RELATIONS

- Stokes extended Newton's observations to multi-dimensions

-replaced "strain" with "strain rate"
---these relationships can not be proved, only verified in practice

STRESS IS A TENSOR

Stress is defined by

- magnitude
- the face on which it acts
- the direction of the force associated with the stress

Temperature is scalar

Velocity is vector

stress is tensor!

(to look at visual examples of this, look on pages 48 and 49)

STOKE'S HYPOTHESIS

$$\lambda = -2/3 * \mu$$

If hydrodynamic pressure conditions are equal to hydrostatic pressure conditions.
Stokes basically made up this relation to μ called λ

And said that there was conservation of mass control volume

The rate at which mass changes inside a fluid element volume is equal to the rate at which it enters

CONTINUITY IN 2D

$$\frac{\partial \rho}{\partial t} + \frac{\partial (\rho u)}{\partial x} + \frac{\partial (\rho v)}{\partial y} = 0 \text{ (page 51)}$$

&**The only assumption is continuum in 2d!!!

In 3d, just add $\frac{\partial (\rho w)}{\partial z}$ at the end

in other words

$$\frac{\partial \rho}{\partial t} + \text{div}(\rho \mathbf{v}) = 0$$

AGAIN, NO ASSUMPTIONS EXCEPT CONTINUUM

If we assume incompressible....

We drop the ρ/bt term and we can bring out the ρ from the div term, and it divides out and we get....

$u/bx + v/by + w/bz = 0$. This is the continuity equation for incompressible flow. IT'S THE SAME AS DILATION!!!! (Page 52)

Derivation of the Navier-Stokes equation....

Need to look at teacher's notes!!! It's on pages 53 and 54

EXACT SOLUTIONS OF NAVIER-STOKES EQUATION

Are rare 😞

There is one situation in which enough of the terms cancel out:

-infinitely long plates

therefore

$\partial/\partial x = 0$ other than pressure term, so $u/bx, v/bx, w/bx = 0$

The pressure derivative $\partial P/\partial x$ is constant, required to drive the flow

Examples of parallel flow

Either a top plate is moving

OR

There is a pressure differential.

Flow between very long parallel plates and flow through a very long pipe

→ parallel flow – no x derivations

→ so

$$\frac{d\rho}{dt} + \rho \frac{du}{dx} + \rho \frac{dv}{dy} = 0$$

If we assume steady flow, then...

There is no variance of density with time so we can get rid of $d\rho/dt$

Then, if we assume incompressible, we can take the ρ term out and it only becomes $du/dx + dv/dy = 0$

If we assume parallel flow then we are only left with dv/dy is equal to zero.

THEREFORE, v is $f(x)$ or v is constant
SO v is constant b/c dv/dx is equal to zero

Apply boundary condition at wall
 $u = v = 0$

If the pressure gradient is zero, $dP/dx = 0$, then this is called Couette flow

if dP/dx is constant this is called a Pouseuille Flow

For dP/dx , look at page 56 and 57 for a derivation and solving of it

you get something like

$$u = Vy/h + 1/\mu * dp/dx * y/2 (y-h)$$

The first term is the component due to the velocity of the upper plate

The second term is the component due to the pressure gradient

For couette flow...

$$u(y) = Vy/h$$

Linear velocity profile

We can also calculate what the shear stress is at the wall

$$T_{xy} = \mu (du/dy + dv/dx) \text{ the } dv/dx \text{ term goes away, to be } T_{xy} = \mu * V/h$$

The coefficient of friction is equal to $T_{wall} / (1/2 * \rho * V^2)$

$$\text{Or....} C_f = 2/Re$$

PLANE POSEUILLE FLOW

Dp/dx is non-zero, $V = 0$ -- there is no velocity of the top plate (this isn't through a pipe)

$$u(y) = 1/\mu * dp/dx * y/2 (y-h) \text{ This is a parabolic profile}$$

When $y-h < 0$, there is a favorable pressure gradient

In a generalized case....

(there is a nice picture of this on page 58)

NOW LET'S LOOK AT FLOW THROUGH A PIPE

Continuity $du/dx + dv/dy = 0$

Cylindrical, so

$$du/dx + 1/r * d(rv)/dr = 0$$

It's infinitely long, so we can assume $d/dx = 0$, except for the pressure gradient term dp/dx

so....

$$d(\rho*u)/dx + 1/r (d(\rho*r*u)/dr = 0$$

take out the first term because it's infinitely long....
simplifies to

$$d(rv)$$

$$/dr = 0 \implies rv \text{ is constant}$$

v at $R = 0$

therefore $v = 0$

1. Flow is steady: $d/dt = 0$

2. flow is $d/d\theta = 0$ (not rotating)

$$u_x = u_\theta = 0$$

$$d/dx = 0$$

therefore you can use this equation:

$$1/r * d/dr * r(du/dr) = 1/\mu * dP/dx$$

Integrate it twice to get

$$u(r) = 1/\mu * dp/dx * r^2/4 + C1 \ln r + C2$$

C1 must equal zero

at $r = R$

left with....

$$u(r) = -1/4\mu * dp/dx [R^2 - r^2]$$

Which is again, a parabolic profile

$dp/dx < 0$ creates a favorable pressure gradient

$dp/dx > 0$ presents an adverse pressure gradient

VOLUME FLOW RATE

Is something that we may run into

It is just the integral from 0 to R of $u(r) * 2 * \pi * r * dr$

Which is....

$$-\pi R^4 / 8 * \mu * dp/dx$$

It's proportional to R^4 not R^2 , which is weird

---Remember, this is all for laminar flow

Average Velocity

$$v = -R^2 / 8\mu * dp/dx$$

This is volume / Area

SKIN FRICTION

$$R/2 * \mu * dp/dx$$

This is just integrating the shear stress

More about Knudsen number

Continuum starts breaking down when Kn is the order of 0.1. If $Kn < 0.1$, you shouldn't assume continuum. If $Kn > 0.1$, then you are safe

Kn equals λ / characteristic length

Some more terms I forgot about:

A.C. (Aircraft Center (maybe?)) is when the pitching moment is independent of α .

The center of pressure is when the pitching moment is 0. This depends on α .

Slip velocityies in fluids

for Liquids...

slip happens when looking at small scale microflows. Such as water in a 30 micro meter microchannel, observed 10% slip

Gasse

-small knudsen number

-micro/nano flow

-rarefied gas (low pressures)

-sometimes hypersonic flows

**basically if you have continuum, you are good assuming no-slip

*****EXACT SOLUTIONS FOR NAVIER-STOKES VECLOCITY PROFIELS

(nice picture on page 7

*****BOUNDARY LAYERS

2D incompressible boundary lyaytr theory

Hydrostatic condition, equivalent to hydrodynamic condition

ν = kinematic viscosity = μ / ρ units are m^2/s

Preliminary remarks about boundary layers

Navier-Stokes equations are 176 years old

-They are non-linear PDE's

-They can't be solved

-superposition of several simple solutions is possible

-as is separation of variables.

Boundary conditions

one on top plate, one on bottom plate for steady flow

Navier-Stokes equations are ELLIPTIC, meaning points in the flow are all coupled to each other. They are computationally expensive to solve!!!

Navier-Stokes is parabolic PDE, which means it depends on one initial condition

N-S for hypersonic is hyperbolic

We need to be able to model the boundary layer
Prandtl (60 years later) decided to try to ignore certain terms
-he did an order of magnitude analysis.

Prandtl divided into 2 regions:

- External region - inviscid assumption
- Boundary Layer - viscous
- characteristic: Very strong gradients: temp velocity

(only relevant for attached boundary layers)

Low alpha still has a boundary layer

Note: a BL is a "shear layer" - a wake is a "free shear layer"
BL rules still apply

ORDER OF MAGNITUDE ANALYSIS (starts on page 70...it's a derivation)The model consists of a very long surface, L , a freestream in the x direction, V , the width of the boundary layer, δ , and that's it.

So, we get
 $u = o(V)$
 $x = O(L)$
 $y = O(\delta)$

Continuity says
 $dy/dx + dv/dy = 0$

$$O(V/L) + O(v/\delta) = 0$$

They must be the same order@!

So you can say that $O(V/L) = O(v/\delta)$

so..

$$v = O(V\delta/L)$$

ADN

$$v/V = O(\delta/L)$$

One of the biggest assumptions is that $d \lll L$

Therefore, $v \lll V$

So, we can look at the u momentum equation.....

$$u \frac{du}{dx} + v \frac{du}{dy} + \frac{1}{\rho} \left(\frac{dp}{dx} \right) = \frac{\mu}{\rho} \left[\frac{d^2u}{dx^2} + \frac{d^2u}{dy^2} \right]$$

The LHS simplifies down to

$$O(VV/L) + O(\delta V)$$

THIS ORDER OF MAGNITUDE ANALYSIS IS ON PAGE 71